

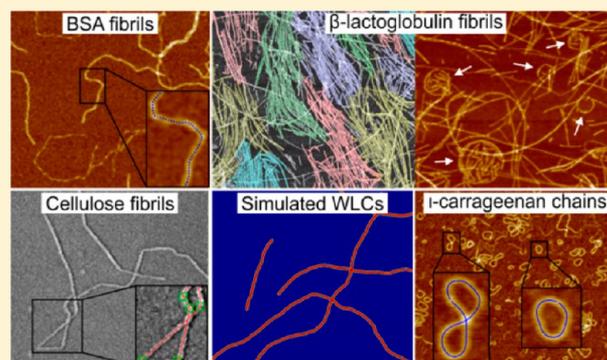
FiberApp: An Open-Source Software for Tracking and Analyzing Polymers, Filaments, Biomacromolecules, and Fibrous Objects

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S Supporting Information

ABSTRACT: Biological semiflexible polymers and filaments such as collagen, fibronectin, actin, microtubules, coiled-coil proteins, DNA, siRNA, amyloid fibrils, etc., are ubiquitous in nature. In biology, these systems have a direct relation to critical processes ranging from the movement of actin or assembly of viruses at cellular interfaces to the growth of amyloid plaques in neurodegenerative diseases. In technology and applied sciences, synthetic macromolecules or fibrous objects such as carbon nanotubes are involved in countless applications. Accessing their intrinsic properties at the single molecule level, such as their molecular conformations or intrinsic stiffness, is central to the understanding of these systems, their properties, and the design of related applications. In this Perspective we introduce FiberApp—a new tracking and analysis software based on a cascade of algorithms describing structural and topological features of objects characterized by a very high length-to-width aspect ratio, generally described as “fiber-like objects”. The program operates on images from any microscopic source (atomic force or transmission electron microscopy, optical, fluorescence, confocal, etc.), acquiring the spatial coordinates of objects by a semiautomated tracking procedure based on A* pathfinding algorithm followed by the application of active contour models and generating virtually any statistical, topological, and graphical output derivable from these coordinates. Demonstrative features of the software include statistical polymer physics analysis of fiber conformations, height, bond and pair correlation functions, mean-squared end-to-end distance and midpoint displacement, 2D order parameter, excess kurtosis, fractal exponent, height profile and its discrete Fourier transform, orientation, length, height, curvature, and kink angle distributions, providing an unprecedented structural description of filamentous synthetic and biological objects.



INTRODUCTION

In modern natural and applied sciences, image processing plays a pivotal role in gaining qualitative and quantitative information from investigated complex systems. A subtype of signal processing, image processing operates on 2D or 3D arrays of numerical values by means of morphological operators, filters of various types, and other standard signal-processing tools. Image processing has a large impact, for example, in neurosciences, via neuron tracing in 3D stacks of fluorescence microscopy images.^{1,2} The extraction of neuron's morphology via accurate localization of soma, axon, and dendrites^{3,4} is a step in reconstruction of neuronal cells' network; that is of great importance for understanding how the brain works.⁵ The image processing is also very helpful in cell biology with segmenting cytoskeletal actin filaments that are involved in cell mobility and division as well as in carrying intracellular transport functions.⁶ From elongation rates over time of these filaments, an estimation of kinetic and equilibrium constants for polymerization and depolymerization reactions allows clarification of detailed actin functionality in cells.^{7,8} DNA is another important object, for which image processing has been traditionally widely used. It has become a classical system to apply polymer physics concepts^{9,10} and to explore statistical

properties and scaling behavior of macromolecules in both linear¹¹ and circular conformations.¹² In the case of amyloid fibrils,^{13–17} understanding the correlations between molecular structure, mechanical properties, and morphology, as obtained from processing atomic force and transmission electron microscopy (AFM, TEM) imaging data, is crucial for elucidating critical steps in the progression of neurodegenerative disorders, such as Alzheimer's, Parkinson's,^{18–20} or sickle cell disease.²¹ Particularly, some recent investigations suggest that disease-specific prion fibrils exhibit higher elastic modulus than prion fibrils without disease specificity.²² By exploiting the stochastic optical reconstruction microscopy, which is essentially a fluorescence microscopy technique based on high-accuracy localization of photoswitchable fluorophores,²³ the image processing can also be used in determining the exchange pathways of monomers within individual supramolecular fibers.²⁴

In order to test the accuracy of tracking procedures and verify the correctness of computational tools, artificial

Received: November 9, 2014

Revised: January 26, 2015

Published: February 10, 2015

generation of synthetic images with curvilinear objects, that follow a certain shape model, can also be used.^{25,26} Simulating images with adjustable size, resolution, level of noise, and system parameters (branching, length, thickness, flexibility, etc.) can greatly help in understanding the limitations of available algorithms and improve their efficiency.²⁷ Hence, gaining access to quantitative and qualitative information via image processing and statistical image analysis is central to a very vast landscape of fields, spanning from in biomedical research to nanotechnology and biomaterials science.

An overarching necessity in all above-mentioned fields is the efficient detection of curvilinear objects in images of various types (2D and 3D), acquired primarily by fluorescence microscopy, confocal laser scanning microscopy, atomic force microscopy, and scanning and transmission electron microscopy (SEM and TEM). This general task has been considered by a vast number of scientists and researchers. The algorithms have been developed based on estimation of the first- and second-order directional derivatives,²⁸ active contour models that work by “energy minimization” of external (image intensities, or gradients) and internal (contour tension and elasticity) forces,²⁹ unbiased detector of curvilinear structures,³⁰ edge and ridge detection based on the optimization of a Canny-like criterion,³¹ or linear feature detection using multiple directional nonmaximum suppression.³² Some of these methods were widely adapted in medical research and neurosurgery for diagnosis of vascular diseases and planning vascular interventions from magnetic resonance angiographic (MRA) images^{33–35} and, particularly, in neuroscience and neurology, where investigation of crucial neuronal functions provides valuable information about the brain activity and possible applications in treatment of neurodegenerative diseases. An additional complication appearing in neuronal tracing is a separate detection of soma and branching points of dendrites, which is normally treated separately from axon tracing.¹ Particularly, an imaging data obtained in these research fields is represented as plane image stacks in 3D that depict fluorescently labeled cultured neurons. The minimal element of such images is called voxel (a combination of words “volume” and “pixel”), which is an image intensity value on a regular grid in three-dimensional space. Generally, the above-mentioned algorithms with a readaptation to 3D can be used independently or in combination with the auxiliary pre- and postprocessing of recorded data. Some of them are based on an object “thinning” (skeletonization) method, primarily via morphological operations of erosion and dilation, that are an iterative removal of voxels from the surface of the object until the last thin line remains.^{4,36–40} However, often they may not be robust enough on noisy images, which may severely disturb the quality of segmentation, and additional image preprocessing by noise filtering is highly required. Alternative methods employ the minimal-cost path calculation (Dijkstra’s or A*⁴¹) between two points defined by the user or, in an automatic way, through an appropriate additional utility algorithm.^{39,42} In a general 2D case, to facilitate the minimal-cost pathfinding algorithm, or as a standalone method, the eigenvalues of Hessian matrix applied to an image are used to compute for each voxel the match to the object and the corresponding eigenvector for its local direction.^{2,43,44} Other methods include voxel scooping⁴⁵ and model-based splines,^{26,44} which were initially introduced as snakes.²⁹

More relevant to macromolecules, fibrillar structures, and polymers, various image processing methods are applied to

quantify bending properties, network geometry, and basic morphological characteristics. Identification of the 3D gel microstructure,⁴⁶ in particular of collagen fibrils forming a matrix,⁴⁷ can be performed via the binarization and thresholding voxels of an image, followed by classification of branches⁴⁸ or by nucleation and local maximum points identification.⁴⁹ The search of centerline for linear, fiber-like objects like actin filaments and microtubules can be implemented in a similar manner,⁵⁰ or by minimal cumulative path cost estimation in tracking of outer tips of microtubules,⁵¹ and quite promising model-based approaches.²⁵ Various methods have been applied for segmenting DNA^{11,52,53} and amyloid fibrils^{54,55} adsorbed on surface and visualized by AFM or TEM. Statistical and scaling properties of these objects can be derived from the tracked contours and from methods for simulation and measurement of amyloid fibril width, length, and height distributions.²⁷

Some macromolecules like DNA or carrageenan polysaccharides may exhibit a circular conformation and can be well described by a closed contour.^{12,56,57} The intrachain bond correlation functions and statistical behavior of these systems are different compared to linear polymers.^{58–60} The tracking approach based on the snakes model²⁹ provides a direct method to address these complications. These closed snakes can be used to obtain spatial description of circular objects in 2D visualized by AFM or EM, with a follow-up analysis of their particular features.

In this context, the development of computational tools for the entire set of linear nonbranched and circular objects is of great importance in nanoscience. There are several examples of programming applications that can be found in the recent literature. Most among them are plug-ins for the Java-based image processing program ImageJ: NeuronJ (neurite tracing and analysis in fluorescence microscopy images),⁴³ Neurite-Tracer (automated quantification of neurite outgrowth),³⁹ JFilament (segmentation and tracking of cytoskeletal filaments),²⁵ and actin filament polymerization–depolymerization tracker.⁷ Furthermore, MATLAB-based or standalone applications are also available: NeuronMetrics (semiautomated processing of cultured neuron images),³⁸ NeuronCyto (neurite outgrowth measurement based on image segmentation with topological dependence),³ V3D-Neuron (3D reconstruction of neurites),⁶¹ DNA Trace,⁶² and Simple Neurite Tracer (reconstruction, visualization and analysis of neuronal processes).⁴² Yet, none of these softwares are fully rooted on polymer physics concepts to analyze filamentous objects, and a comprehensive polymeric description of filamentous objects is yet to be provided within a single program application, which calls for immediate developments.

Over the past years an increase in acquisition modes of experimental visualization techniques and a growth of their quality and resolution have been remarkably rapid. It is noticeable, especially, in AFM imaging with more accurate cantilevers and possibilities to scan very-high-resolution images of the size 5120 × 5120 pixels. With the parallel progress in computer speed and memory volumes, more precise image analysis tools and improved data handling have become highly desirable. In this Perspective we present a solution to above-mentioned application problems raised in nanoscience—a project with the open-source code called FiberApp, written in MATLAB programming language, in which cascading algorithms based on statistical polymer physics concepts have been implemented to yield the most accurate conformational,

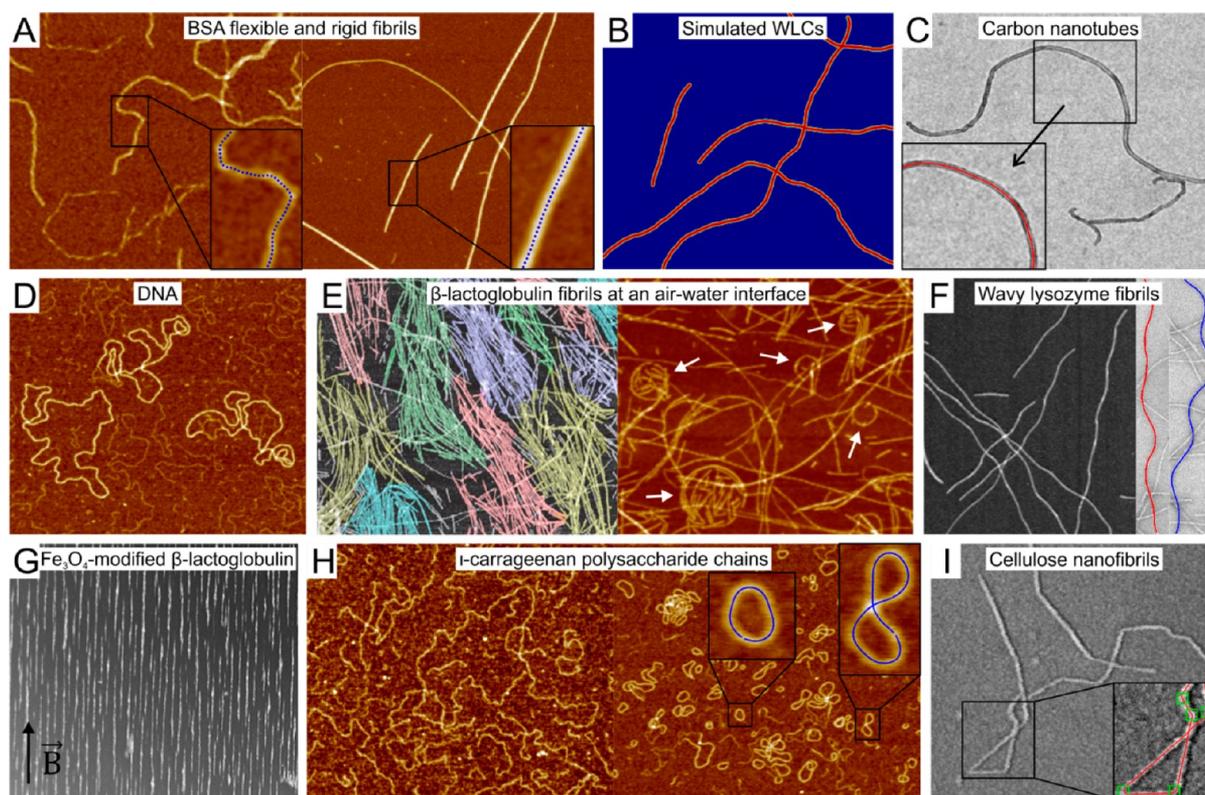


Figure 1. A few illustrative examples of systems that can be analyzed by FiberApp software. (A) Bovine serum albumin (BSA) flexible (left) and rigid (right) amyloid fibrils.^{63,64} The inset images depict magnified areas of AFM images with point positioning along tracked contours with low and high intrinsic stiffness, respectively. (B) Simulated worm-like chain fibrils, generated directly in FiberApp software, providing a theoretical benchmark model against real systems.⁶⁶ (C) Functionalized multiwalled carbon nanotubes (the figure is adapted from Li et al.⁶⁷). The inset image reflects the feasibility of tracking for hollow objects. (D) Circular DNA adsorbed on APTES-modified mica (with carrageenan fibrils in the background). Image courtesy of Larissa Schefer. (E) β -Lactoglobulin fibrils that form 2D liquid crystalline domains (left)⁶⁵ and circles (right), pointed by white arrows, at liquid interfaces (figure adapted from Jordens et al.⁶⁶). (F) Wavy lysozyme amyloid fibrils with subsistence-length complex scaling behavior (figure adapted from Lara et al.⁶⁹). (G) Nanoclusters of Fe_3O_4 nanoparticles with β -lactoglobulin fibrils that can align in the presence of a magnetic field (figure adapted from Bolisetty et al.⁷⁰). (H) Linear *t*-carrageenan polysaccharide chains, prior (left) and after addition of salt (right), forming secondary structures and looped conformation.⁵⁷ Image courtesy of Larissa Schefer. The inset images demonstrate the possibility of closed-contour tracking. (I) TEMPO-oxidized wood cellulose nanofibrils with areas of different intrinsic stiffness (kinks), originating from the harsh mechanical treatment during sample preparation.⁷¹ In the inset image it is demonstrated a concept of tracking with special masks that define contour segments with high curvature and low intrinsic stiffness.

topological, and structural analysis available to date to describe fibrous and filamentous biological and synthetic objects. The input images, on which FiberApp works, may have been generated by any microscopy technique, essentially any TIF image, although most of the examples which will be presented are based on high-resolution atomic force microscopy and scanning and transmission electron microscopy imaging. We employ the most suitable and promising tracking method based on active contour models, reinforced by the additional utility A* pathfinding algorithm for a faster semiautomatic segmentation process, allowing the tracking of thousands of objects in a short time. The potential to track linear and circular fiber-like objects and, independently, the possibility to define segments of heterogeneous stiffness along their contours using mask elements are only some of the key features that broaden the scope of possible applications of the program to a context much wider than that of classic statistical polymer physics. Moreover, the ability of FiberApp to simulate “in silico” statistical data and artificial images of fiber-like objects serves as a powerful tool for prediction and rationalization of experimental results.

DESCRIPTION OF THE SOFTWARE

Typical Systems of Interest Handled by FiberApp Software. The concept of representing fiber-like objects solely by their midline contours, tracked in microscopy images, is applicable to a broad variety of systems in nanoscience. A few types of such objects and their features, which have previously been investigated with the help of FiberApp software, are exemplified in Figure 1. Figure 1A shows the application of FiberApp to bovine serum albumin (BSA) flexible (Figure 1A, left) and rigid (Figure 1A, right) amyloid fibrils, which are formed upon incubating the protein solution at an elevated temperature (90 °C) and in an acidic environment (pH 2) for several days. This analysis allowed the identification of six distinct polymorphic types of BSA fibrils, which coexist with four types of intermediates. Furthermore, the handedness inversion, upon a level change in their hierarchical organization, was resolved.^{63,64} Another example is the statistical analysis of simulated worm-like chain (WLC) contours with known and predefined structural parameters (Figure 1B).⁶⁵ These artificial images of simulated WLC fibrils, with the same resolution as the AFM imaging data subjected to the tracking procedure, can help to rule out possible artifacts that originate from tracking

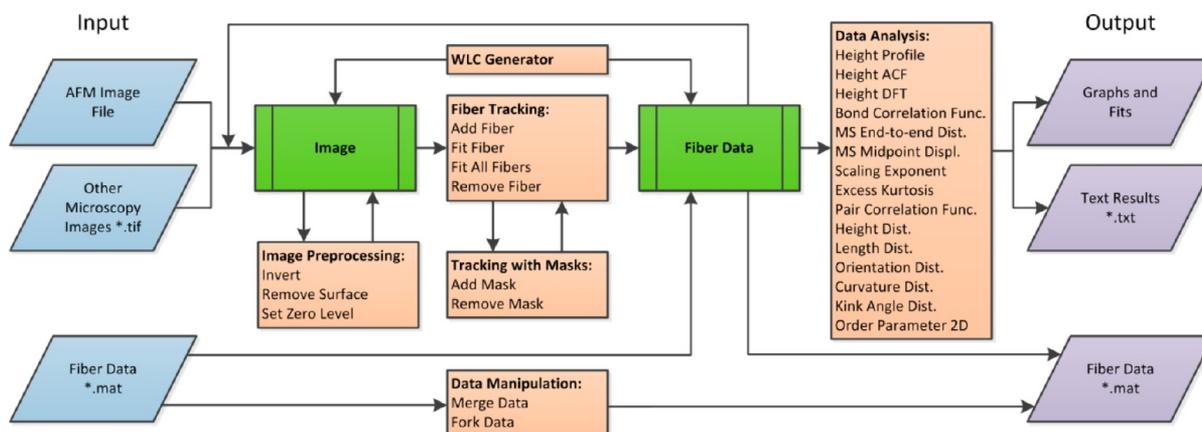


Figure 2. General scheme of the workflow of FiberApp software. The possible input data types are given at the left side of the figure (blue parallelograms). The main internal data types of the program “Image” and “Fiber Data” (green rectangles) define a framework of cyclic tracking procedure, where fiber data can be collected from, potentially, several images. The program’s core functionalities are organized according to their specifications and the areas of implementation (orange rectangles). The functions for the fiber data analysis result in forms of output given at the right side of the figure (violet parallelograms).

algorithms and, hence, be useful for analyzing the real fibril images.⁶⁶ Rigidity and mechanical properties of functionalized multiwalled carbon nanotubes with sulfonic groups are also accessible via FiberApp (Figure 1C).⁶⁷ Topological characteristics of DNA,⁵⁷ which is also an extensively used system for the application of polymer physics concepts in general, as handled within FiberApp, are shown in Figure 1D.^{10,68} Another example, which consists of the liquid crystalline alignment of β -lactoglobulin fibrils at liquid interfaces, where the coexistence of two-dimensional isotropic and nematic phases can occur, is resolved and quantified by a length-scale-dependent order parameter S_{2D} calculated on all tracked contours (Figure 1E, left).⁶⁵ The same polar β -lactoglobulin fibrils at the air–water interface do not possess the characteristic Gaussian curvature distribution for 2D WLC but instead have an excess of large curvature values (fat-tailed curvature distribution), which can lead to ring-like fibril conformations (Figure 1E, right).⁶⁶ In the case of wavy lysozyme fibrils, prepared by incubation of the protein solution at pH 2 and 60 °C for several days, polymer physics concepts were applied via FiberApp on tracked contours in order to study the subpersistence-length complex scaling behavior and the periodicity in height profiles (Figure 1F).⁶⁹ Another class of objects handled by the software is a system of magnetic-responsive Fe_3O_4 nanoparticle-modified protein spherical nanoclusters that can align in the presence of a magnetic field (Figure 1G).⁷⁰ The system of *t*-carrageenan polysaccharide chains in the random coil conformation (Figure 1H, left) that undergoes a coil–helix transition at high ionic strength coupled with an increase in thickness and rigidity of polymers is another example of biomacromolecules which can be handled by the software (Figure 1H, right). Along with the formation of the secondary structure, *t*-carrageenan chains also form looped structures that coexist with linear polymer chains.⁵⁷ TEMPO-oxidized wood cellulose nanofibrils with kinks provide an example of heterogeneous stiffness along the contours, which causes an additional complication for the tracking procedure. FiberApp solves this problem by employing masks that define kink areas, so that the contour affinity to bend depends on whether the contour segment is inside or outside the mask (Figure 1I). This is essential for the tracking algorithm to correctly follow the nanocellulose fibrils⁷¹ or DNA with protein-induced bending.^{72,73} There is also the possibility

to apply more sophisticated theories for chains with bends or sections of a different flexibility.⁶⁸ In the case of nanocellulose fibrils, the detailed statistical investigation of the kink angle distribution provided convincing evidence that the commonly accepted model of the structure as built of alternating hard crystalline and soft amorphous regions of cellulose chains is not appropriate and supported a process-induced kink formation generated during the preparation treatment.⁷¹ Other examples of supramolecular fibrous aggregates, not shown in Figure 1, with periodic height changes accessible via FiberApp, are offered by self-assembled glycyrrhizic acid forming twisted ribbons⁷⁴ or ILQINS hexapeptides forming helical ribbons.⁷⁵

Description of the Workflow with FiberApp Software.

Together with a brief description of A* pathfinding and active contour models based tracking procedures, modified and adapted to the needs of the software, we provide a complete workflow scheme for processing AFM and other types of microscopy images, as shown in Figure 2. The code is optimized for working with modern superhigh-resolution microscopy images (5120×5120 pixels and virtually more) minimizing the number of errors upon tracking of the objects. Functions in the “Image Preprocessing” group can be used prior to fiber tracking in order to remove uneven illumination from the image and allow correct extraction of z coordinates.

There are two major steps over the general workflow: “Fiber Tracking” and “Data Analysis”. With the first step, each fiber-like object in the image is fully characterized by its contour, which is a sequence of points ordered along the fiber’s middle line from one end to the other. These contours have a constant distance between the point projections on the image plane, the step size Δs , which has a typical value of about one to several pixels size of the image. Thus, the coordinates x_i , y_i (in the image plane), and z_i (height) of all the points covered by the contour provide a full spatial description of the fiber’s central line.

The initial positioning of a contour is manually defined with the help of a modified A* pathfinding algorithm.⁴¹ For that, the image is considered to be a weighted graph of “moving costs”, where image pixels are represented by graph vertices. The “cost of movement” between two adjacent pixels is inversely proportional to the starting pixel intensity, and this value is stored as a graph edge between two corresponding graph

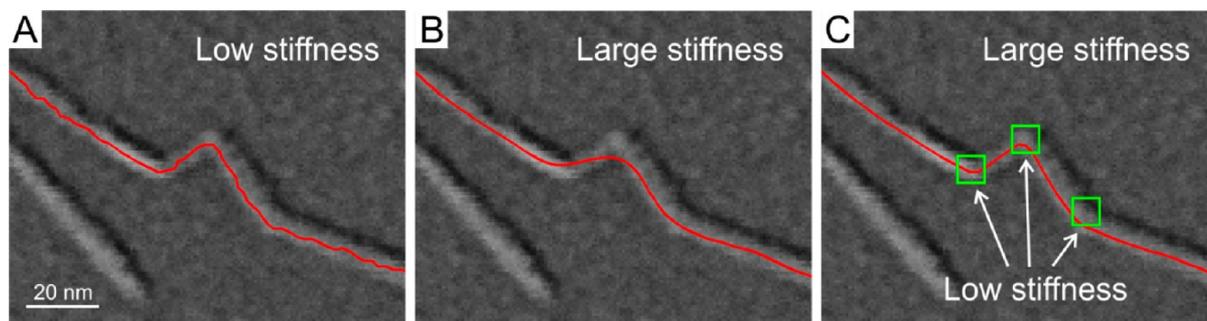


Figure 3. Cryo-SEM images with tracked contours of nanocellulose fibrils. (A) A low stiffness parameter was used to perform tracking, which leads to an abrupt contour with undesirable fluctuations along it due to noise in the image. (B) The opposite case, a large stiffness parameter, suppresses the noise influence but does not allow an accurate fit in the vicinities of kinks. (C) The special masks (green squares) can be used to provide the contour with a heterogeneous stiffness, by making contour segments inside the masks softer (or stiffer) in comparison to the outer parts.

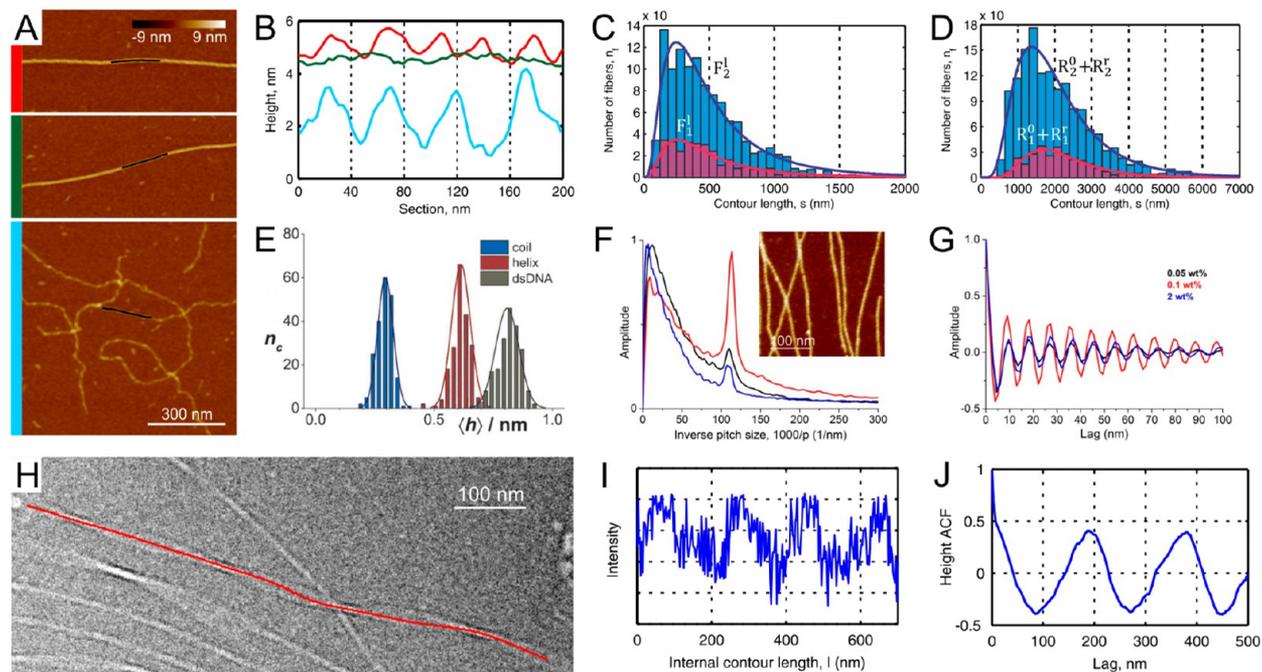


Figure 4. (A) AFM images with BSA fibrils of type F_1^1 (bottom), R_1^0 (middle), and R_1^1 (top). The scale bar and the color bar apply to all AFM images. (B) Height profiles of longitudinal sections along fibrils with respective color identifications from the panel A. The figures in panels A and B are readapted from Usov et al.⁶³ (C) Length distributions of flexible thin F_1^1 and flexible thick F_2^1 classes of BSA fibrils acquired after 40 h of incubation. (D) Length distributions of rigid thin $R_1^0 + R_1^1$ and rigid thick $R_2^0 + R_2^1$ classes of BSA fibrils acquired after 100 h of incubation. The solid lines represent the best-fits for histograms with a log-normal distribution function. The figures in panels C and D are readapted from Usov et al.⁶⁴ (E) Average height distributions of *t*-carrageenan polysaccharide chains in random coil and ordered helical conformations and double-stranded DNA. The figure in panel E is reproduced from Schefer et al.⁵⁷ (F, G) Periodicity analysis and the pitch size estimation based on Fourier transform and autocorrelation function of the height profiles of supramolecular aggregates at different concentrations. The inset figure depicts an AFM image of analyzed fibrils. The figures in panels F and G are readapted from Saha et al.⁷⁴ (H) Cryo-TEM image of twisted lysozyme amyloid fibrils and a tracked contour shown as a red line. Image courtesy of Stephan Handschin. (I) Noisy pixel intensity profile along the contour. (J) Height autocorrelation function with clear peaks defining the periodicity of lysozyme amyloid fibrils.

vertices. The algorithm then heuristically searches for the “lowest cost path” between two points that are manually defined by a user. Hence, the path will preferably go along a bright ridge of a fiber-like object in the image, which is a good first approximation of the object’s midline position.

The contour is, afterward, iteratively and automatically deformed according to the active contour model.^{25,29} These contours are parametric curves that adapt to the image features in order to minimize the “total contour energy”, which is expressed as the sum of an “external energy”—increasing with the positional offset of the contour from the brightest pixels of

the image—and an “internal energy”—increasing with the growth of the total curvature along the contour.

Functions in “Tracking with Masks” provide a possibility to use special mask elements that can be placed to define areas of the contour with different bending properties (Figure 3). For example, the curvature-penalizing parameter can be locally reduced, so that the contour can more easily bend there and exhibit a kink. In other words, it allows contours to have a heterogeneous stiffness (e.g., low stiffness inside mask areas, large stiffness outside). In this example, the masks are used to force the contour to bend in the kink areas and also define an angle of contour deviation (kink angle) for further analysis. The

resulting coordinates of contours and masks after the tracking procedure can then be stored in a separate project file and provide the basis for further data analysis.

Furthermore, the tool “WLC Generator” is used for creating both synthetic contour coordinates that follow WLC statistics and images with fibrils scattered on them, while functions in the “Data Manipulation” set provide a simple solution for combining and splitting statistics from different images or fiber types. In what follows, we describe in detail the second step of the workflow with FiberApp—the basics of fiber data analysis that can be performed and typical results output that can be generated. The modulated internal architecture of FiberApp allows incorporating new processing methods with ease, thus making it a flexible platform for future developments. The technical details of processing algorithms, together with specifications of A* pathfinding and active contour models based tracking procedures, are discussed in detail in the Supporting Information.

SPECIFICATIONS AND OUTPUTS OF THE IMPLEMENTED FIBER DATA PROCESSING METHODS

Basic Morphological Parameters of Fiber-like Objects and Periodicity along Their Contours. The methods in this section provide a primary morphological characterization of the objects accessible by the software and include the height profile, the height and length distributions, the height autocorrelation function, and the height discrete Fourier transform analysis.

The height profile is a simple test to briefly estimate the average height and check the presence of repeating patterns along the contour. It also can serve in distinguishing between different families of fiber-like objects. For example, Figure 4B represents the longitudinal sections of the BSA fibrils displayed in Figure 4A, by which three distinct types of thin BSA fibrils (flexible left-handed F_1^1 , rigid without periodicity R_1^0 , and rigid right-handed R_1^1) based on their average height, existence of periodical structure along their contours and visual appearance, can be determined.⁶³ This method can be used to plot height profiles not only along curved objects but also along linear cross sections.

The length distribution can provide information about nucleation and growth rates of fibrils, when distributions of samples, taken at different incubation times, are compared. Furthermore, in the system of BSA fibrils the statistical results indicate that the conversion rate of the polymorphic transformation $F_1^1 \rightarrow R_1^0 + R_1^1$ is slightly lower than that of the transformation $F_2^1 \rightarrow R_2^0 + R_2^1$. This arises from the fact that the total length ratio of flexible classes $L(F_1^1)/L(F_2^1) \cong 0.27$ at 40 h of incubation is higher than the ratio of rigid classes $L(R_1^0 + R_1^1)/L(R_2^0 + R_2^1) \cong 0.2$ at 100 h.⁶⁴ The shape of length distribution for amyloid, and cellulose fibrils,⁷¹ as well as many other parameters in biology⁷⁶ fits well to the log-normal probability density function $f(L)$:

$$f(L) = \frac{A}{L\sigma\sqrt{2\pi}} e^{-(\ln L - \mu)^2 / 2\sigma^2} \quad (1)$$

where L is the total length, μ and σ are the mean value and the standard deviation of the length natural logarithm, respectively, and A is a normalizing constant (Figure 4C,D).

By plotting histograms of the average height values, extracted mainly from AFM images, it then becomes possible to analyze and discriminate between different families of fiber-like objects.

An example is shown in Figure 4E, where the height distributions are plotted for *t*-carrageenan polysaccharide chains in random coil and ordered helical conformations together with the double-stranded DNA. From this accurate analysis, it was possible to resolve a long-standing question whether the ordered helical conformation of carrageenans exists as a single or double helix. The final conclusion that can be drawn states that thicker strands consist only of a single polymer chain.⁵⁷

Discrete Fourier transform (DFT) of height profiles provides a way to estimate a pitch size of objects with periodical height variations as in twisted or helical ribbons. Fast Fourier transform (FFT) is a rapid method to calculate DFT and has the form

$$A_k = \sum_{n=1}^N h_n e^{-2\pi i(k-1)[(n-1)/N]} \quad k = 1, \dots, N \quad (2)$$

where h_n is the height of the n th point along the contour, N is taken as a maximal number of points in a contour among all fiber-like objects, and i is the imaginary unit. Amplitudes A_k of the DFT describes the presence of a frequency component k in the height profile. One clear peak in the graph is the signature of a distinct periodical structure, whose value in nanometers can be estimated via $1000/(\text{peak position})$.

The second way to study periodical patterns is to calculate the height autocorrelation function R_{xx} . The discrete autocorrelation function has the general form

$$R_{xx}(m) = \frac{1}{N-m} \sum_{n=0}^{N-1-m} x_n x_{n+m} \quad (3)$$

where m is the lag, N is the total number of points along a contour, and x identifies either the discrete height h_n or twist angle θ_n profile. The twist (helix) angle profile is obtained via the height profile's maxima and minima peaks position. The angle values span linearly from 0 to π within the distance between maxima and the consequent minima height peaks and from π to 2π until the next maxima to complete the period. Hence, for sake of simplicity, the twist angle and height profiles have same periodicities (note that often it is assumed that a full period for the angle should have a doubled value in comparison to the height). A presence of periodicity in the resulting curve is a signature of the periodical pattern in height/twist angle profiles with the same period.

Both the above-mentioned methods were applied to yield consistent values for the periodicity ($p = 8-9$ nm) of self-assembled fibrillar supramolecular structures of glycyrrhizic acid (Figure 4F,G).⁷⁴ We note here that the information on the pitch size is also essential for estimation of the mechanical properties of fiber-like objects in helical ribbons conformation.⁷⁷ Moreover, Figure 4H shows a Cryo-TEM image of twisted lysozyme amyloid fibrils with a tracked contour (red line). Despite the fact that the height profile along the contour is very noisy (Figure 4I), the Height Autocorrelation Function results in the clear peaks that define the periodicity of these lysozyme amyloid fibrils (Figure 4J) showing a great robustness of the approach.

Persistence Length Estimation Methods. The persistence length is a basic property of polymers, which quantifies their rigidity and is formally defined via the bond correlation function in 3D as the length over which angular correlations in the tangent direction decrease by e times.⁷⁸ There are several approaches of the persistence length estimation for WLC

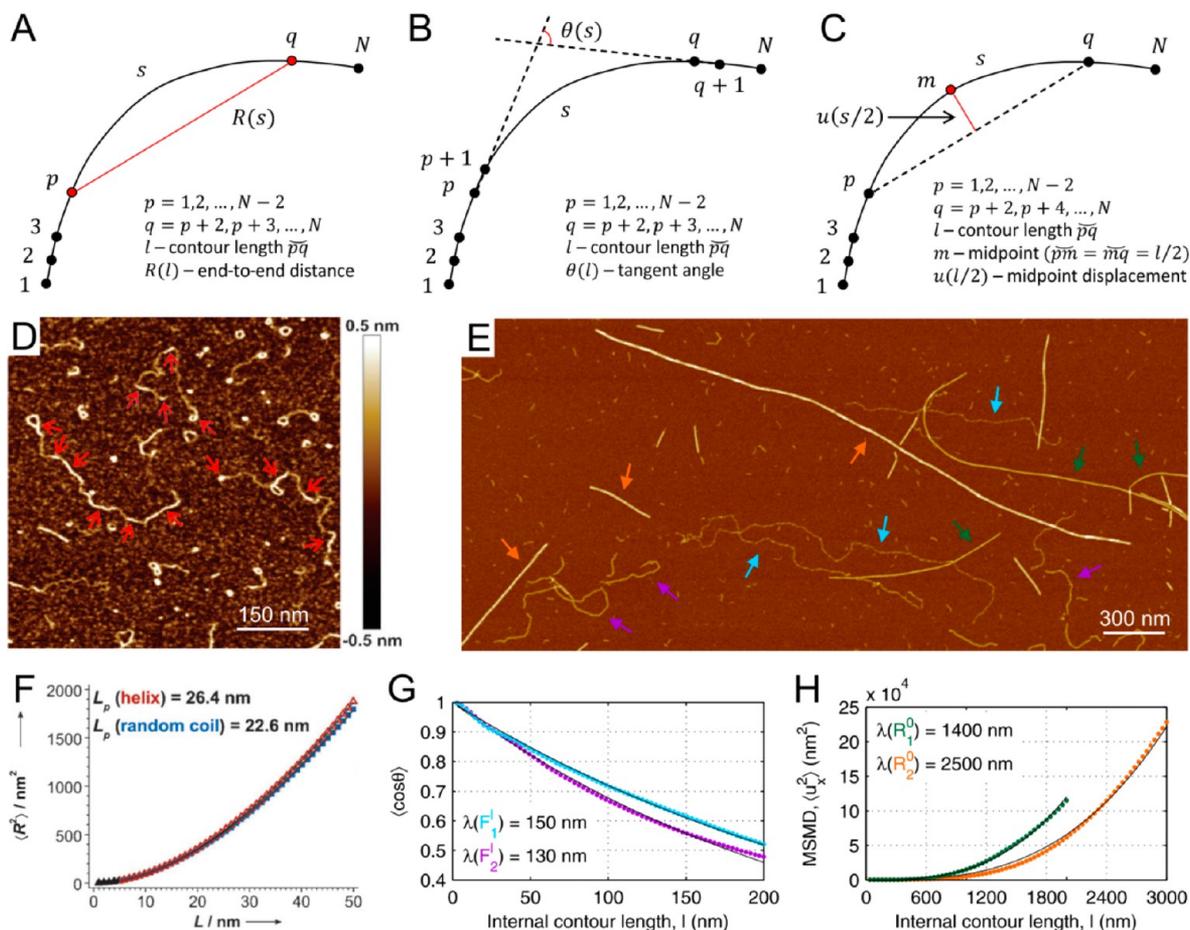


Figure 5. Overview of the persistence length estimation methods for discrete contours that are employed in FiberApp and their practical applications. Schematic representations of methods: (A) mean-squared end-to-end distance (MSED) versus internal contour length; (B) bond correlation function (BCF); (C) mean-squared midpoint displacement (MSMD) versus internal contour length. (D) AFM image of *t*-carrageenan polysaccharide chains in random coil and ordered helical conformations. The red arrows point toward the thick helix segments. (E) AFM image of BSA fibrils on mica. The arrows of the same color point to fibrils of a certain type: F_1^1 (cyan), F_2^1 (violet), R_1^0 (green), and R_2^0 (orange). (F) Plot of MSED versus internal contour length for separate helix and random coil segments of the *t*-carrageenan polysaccharides and values of corresponding persistence lengths extracted from fits. (G) Plot of BCF for BSA fibrils of types F_1^1 and F_2^1 with resulting persistence lengths. (H) Plot of MSMD versus internal contour length for BSA fibrils of types R_1^0 and R_2^0 with resulting persistence lengths. The figures in panels D and F are readapted from Schefer et al.⁵⁷ The graphs in panels G and H are readapted from Usov et al.⁶³

discrete contours, consisting of N points and, accordingly, $N - 1$ segments (Figure 5A–C). However, there is evidence that some systems like DNA at short length scale^{52,79} or simulated bottle-brush polymers when the chain length of the macromolecules tends to infinity⁸⁰ do not necessarily follow the standard statistics of WLC but can be well described by the alternative generalized theory of semiflexible polymers. Here, we assume a WLC behavior for investigated systems of fiber-like objects at all length scales.

One of the most practical and widely used methods for the persistence length estimation is to calculate the mean-squared end-to-end distance (MSED) between contour segments (Figure 5A). This characteristic for a WLC model in 2D has the following theoretical dependence: $\langle R^2 \rangle = 4\lambda[l - 2\lambda(1 - e^{-l/2\lambda})]$,⁹ where λ is the persistence length and R is the direct distance between any pair of segments along a contour separated by an arc length l . Employing this approach, a subtle increase in persistence length of *t*-carrageenan chains upon a coil–helix transition from 22.6 to 26.4 nm could be detected, indicating an increase in rigidity of segments in helical

conformation (Figure 5D,F) and yielding a valuable direct indication of the coil–helix transition on a single chain level.⁵⁷

The bond correlation function (BCF) is the most general way to evaluate the persistence length (Figure 5B). For WLC in 2D it corresponds to $\langle \cos \theta \rangle = e^{-l/2\lambda}$,⁸¹ where θ is the angle between tangent directions of any two segments along a fibril contour separated by an arc length l . A different method that can be successfully applied only to very stiff fiber-like objects ($l < \lambda$) is the mean-squared midpoint displacement (MSMD) as depicted in Figure 5C. The equation, describing the behavior of a midpoint deviation has the form $\langle u_x^2 \rangle = l^3/48\lambda$,^{21,55} where $\langle u_x^2 \rangle$ is the mean-squared midpoint displacement between any pair of segments along a contour, separated by an arc length l . This expression is derived with an assumption that these deviations are small in comparison to the corresponding arc lengths ($|u_x| \ll l$). Both methods were successfully applied for the persistence length estimation of BSA fibrils. The flexible fibrils were analyzed by means of BCF (Figure 5E,G), while rigid fibrils were via the MSMD method (Figure 5E,H). The obtained results are consistent with MSED analysis for all classes (λ_{MSED} is equal to 160 nm for F_1^1 , 150 nm for F_2^1 , 1280

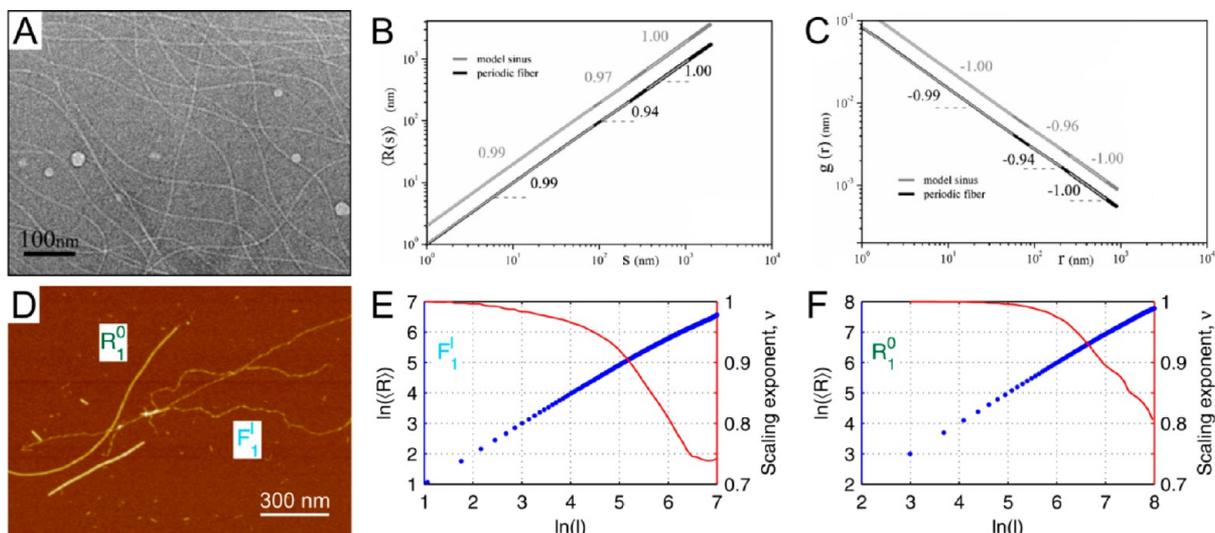


Figure 6. (A) Cryo-TEM image of lysozyme amyloid fibrils. (B) Averaged end-to-end distance versus the internal contour length plot with the scaling exponent ν values for sinus model and lysozyme fibrils at three different length scales. (C) Pair correlation function plot with corresponding values for the scaling exponent α . The figures in panels A, B, and C are readapted from Lara et al.⁶⁹ (D) AFM image of BSA amyloid fibrils of types F_1^1 and R_1^0 . (E, F) Averaged end-to-end distance versus the internal contour length plots (blue axis) and scaling exponent propagations (red axis) for BSA fibrils of types (E) F_1^1 and (F) R_1^0 . The figures in panels E and F are readapted from Usov et al.⁶³

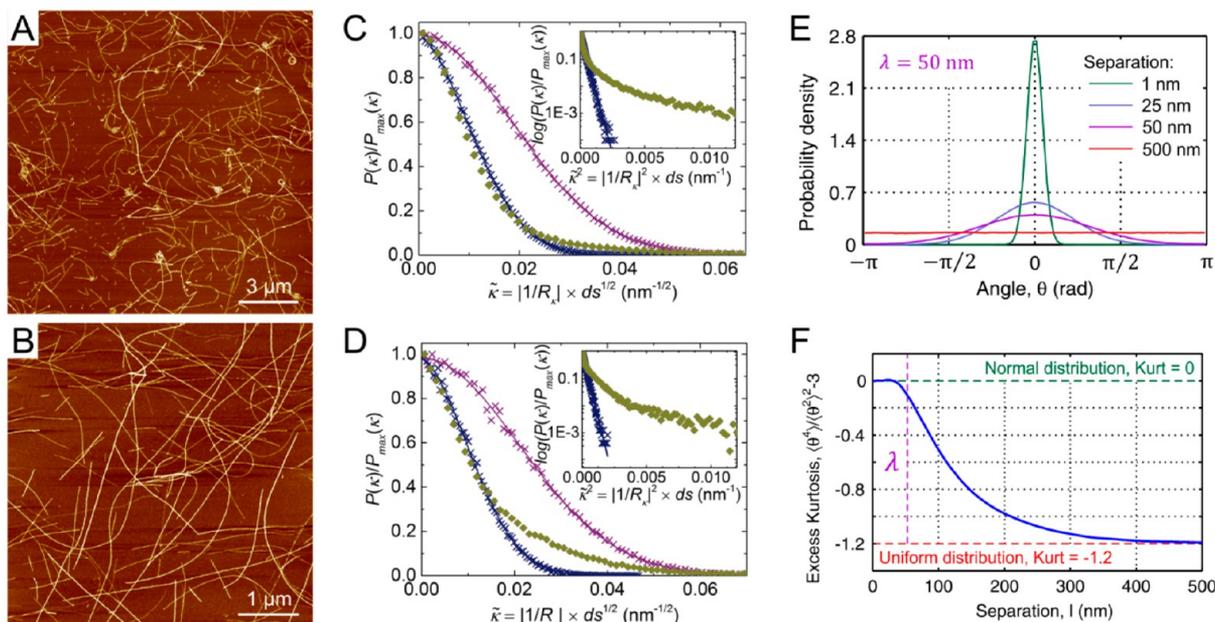


Figure 7. AFM image of β -lactoglobulin fibrils (A) at the air–water interface with ring configurations and (B) deposited onto mica from the bulk solution. (C, D) Normalized probability distributions of absolute curvature values extracted from the generated WLC data (violet), the generated WLC data after applying the tracking procedure (blue), real fibrils from panels A and B, respectively (green). In both scenarios real fibrils exhibit an excess of curvature upon being adsorbed to a surface, which is clearly visible in the inset images. The figures in panels A, B, C, and D are reproduced from Jordens et al.⁶⁶ (E) Probability density functions of the angles between segments of discrete WLCs with identical contour length $L = 700$ nm, persistence length $\lambda = 50$ nm, and discretization step size $\Delta s = 1$ nm. The segments are separated by a distance 1 nm (green), 25 nm (blue), 50 nm (violet), and 500 nm (red). (F) Excess kurtosis as a function of the separation distance between segments of WLCs.

nm for R_1^0 , and 2550 nm for R_2^0)⁶³ and give access to bending rigidities D of the corresponding fibrils as a characteristic of mechanical properties through the expression $D = \lambda k_B T$.⁸²

Scaling Behavior and Pair Correlation Function. Based on polymer physics concepts, the average end-to-end distance is a power-law function of the internal contour length l and scales according to the relation $\langle R(l) \rangle \sim l^\nu$,⁸³ where ν is a scaling exponent. When l is small in comparison to the persistence length ($l < \lambda$), the scaling exponent approaches the

value 1, indicating rod-like statistical behavior. For the contour lengths above the persistence length $l \gg \lambda$ the expected scaling exponent might vary, depending on temperature, quality of solvent, excluded volume interaction, and other parameters, and for the important case of self-avoiding random walk (SAW) chains in a good solvent this is equal to 3/4 in 2D and 3/5 in 3D.¹⁰

The 2D pair-correlation function $g(r)$ is defined as $g(r) = \langle s(r) \rangle / \pi r^2$,⁸³ where $\langle s(r) \rangle$ is the length of the contour segment

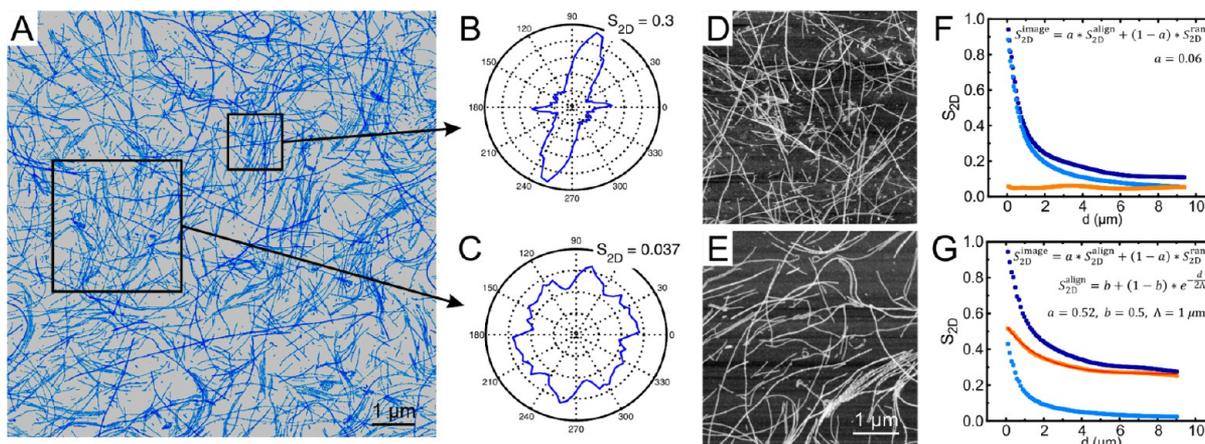


Figure 8. (A) AFM image of β -lactoglobulin fibrils adsorbed at the air–water interface. (B, C) Orientation distribution of contour segments within areas highlighted in panel A. The 2D order parameter within fibril nematic domains is relatively large (panel B, $S_{2D} = 0.3$) in comparison to the corresponding value in the isotropic region (panel C, $S_{2D} = 0.037$). (D, E) AFM images of β -lactoglobulin fibrils adsorbed at the air–water interface after a short period (10 min) and a long period (60 min) of adsorption times, respectively. (F, G) Decompositions of 2D order parameter S_{2D}^{image} of the images from the panels D and G, respectively, into S_{2D}^{align} and S_{2D}^{rand} with weight factors a and b . The images in panels A, D, E, F, and G are readapted from Jordens et al.⁶⁵

incorporated in a circle of radius r , where the average is taken by centering the circles at the contour points with all possible positions of the circle. In the rod-like limit ($l < \lambda$) the pair correlation function behaves as $g(r) \sim r^\alpha$, which together with the approximation $2r \approx \langle R(s) \rangle \sim s^\nu$ results in a relationship between two exponents $\alpha = 1/\nu - 2$.⁶⁹

Using FiberApp, the estimation of exponents for sinusoidal lysozyme amyloid fibrils below the persistence length ($2.5 \mu\text{m}$) was carried out by both approaches (Figure 6A–C), with the observable progressive changes in the values of ν and α at multiple length scales.⁶⁹

An alternative way of presenting scaling exponents, based on the analysis of two BSA fibril classes F_1^1 and R_1^0 (Figure 6D), is illustrated in Figure 6E,F. A local linear fit of the log–log plot with the end-to-end distance versus the internal contour length dependence reflects the scaling exponent at this particular position. The resulting curve is relatively smooth, and in such a case, it can accurately express the evolution of the scaling exponent value with contour length. For different systems the scaling exponent can indicate whether the fiber-like objects are in equilibrated or “trapped” configurations due to the interaction with a substrate.^{9,84} For example, in the case of the flexible BSA fibrils F_1^1 , the scaling exponent reached the plateau of $3/4$ expected for SAW in 2D, indicating the equilibrated state on mica. For the rigid fibrils R_1^0 , on the other hand, the processing length is not long enough compared to the persistence length to draw any conclusion and the scaling exponent does not attain any plateau value, i.e., is not stationary around the $3/4$ expected value.⁶³

Curvature Distribution and Excess Kurtosis. The analysis involving the curvature estimation for polar and chiral fibrils can shed light into important aspects of interaction with a surface upon their adsorption. The probability density function $p(\kappa)$ for the curvature κ of WLC in 2D is equal to $p(\kappa) = (2\lambda\Delta s/\pi)^{1/2}e^{-\lambda\Delta s\kappa^2}$ and is described by a Gaussian distribution of κ .^{25,85} It was shown that the fibrils exposed to an inhomogeneous environment, such as an air–water interface or a mica substrate (Figure 7A,B), do not possess the same Gaussian distribution as WLCs. In all cases, there is an evidence of an excess of curvature, which can also lead to ring-like

conformations of these fibrils (Figure 7A). This spontaneous curvature can be easily detected in the fat tails of the real fibrils curvature distributions in comparison to the synthetic WLCs data, which were subjected to the same tracking procedure in the corresponding simulated images, but still retained the Gaussian shape of curvature distribution (Figure 7C,D). The simulated images of WLCs contours were generated using parameters from the corresponding AFM images, by duly taking into account the finite tip size effect.^{86,87}

The equilibration of the fiber-like objects adsorbed on a substrate can also be verified using the excess kurtosis of the angle θ between consecutive segments of a tracked discrete contour.^{11,53} The excess kurtosis is defined as $K = \langle \theta^4 \rangle / \langle \theta^2 \rangle^2 - 3$, and if the chains are fully equilibrated on a substrate in 2D, then the distribution of θ is Gaussian and $K = 0$, which is a feature of the WLC behavior.⁸⁸ However, it is only possible to apply this method on separation distances below the persistence length ($l < \lambda$) because upon increasing the distance between segments the excess kurtosis value starts to decrease down to $K = -1.2$, and the probability density function of the angle θ transforms from Gaussian into uniform distribution shape (Figure 7E). The corresponding evolution of the excess kurtosis is depicted in Figure 7F for a population of 7000 generated WLC with identical contour length $L = 700$ nm, persistence length $\lambda = 50$ nm, and step size $\Delta s = 1$ nm.

Orientation Distribution and 2D Order Parameter. The orientation distribution and the length-scale-dependent 2D order parameter can be used for quantifying the alignment of fiber-like objects in different systems and under various conditions. Figure 8A shows the AFM image of β -lactoglobulin fibrils adsorbed at the air–water interface with the coexistence of strongly aligned nematic domains and completely randomly oriented fibrils, which form an isotropic phase in analogy to liquid crystals. The orientation distribution of segments belonging to all tracked contours—inside box-size-dependent domains—exhibits a clear peak along one direction (Figure 8B), while randomly oriented fibrils possess a uniform distribution in all possible directions (Figure 8C). It is possible to quantify the order of the assembly by assigning a particular number for each distribution, i.e., via a 2D order parameter.

This is defined as $S_{2D} = 2\langle \cos^2\theta_n \rangle - 1$, where θ is the angle between the n th segment and the local director in the chosen area. The values of the 2D order parameter for the above-mentioned areas are very different: 0.3 and 0.037, respectively (Figure 8B,C). This measure was used to quantify, for example, the alignment of the magnetic-responsive clusters containing hybrids of Fe₃O₄ nanoparticles with β -lactoglobulin amyloid fibrils upon application of a magnetic field (Figure 1G).⁷⁰

In the case of β -lactoglobulin fibrils adsorbed at the air–water interface we introduced an additional feature, which is the length scale dependence of the 2D order parameter, i.e., $S_{2D}(d)$. For calculating this object, we divided the whole image into square blocks of a certain size d . Calculating and averaging S_{2D} values for all blocks results in one mean number, which is parametric with d , yielding the length scale dependent $S_{2D}(d)$.⁶⁵

This function $S_{2D}(d)$ is further expressed as the sum of the weighted components $S_{2D}^{\text{align}}(d)$ and $S_{2D}^{\text{rand}}(d)$, corresponding to the alignment of the nematic and isotropic components, respectively⁶⁵

$$S_{2D}^{\text{image}}(d) = aS_{2D}^{\text{align}}(d) + (1 - a)S_{2D}^{\text{rand}}(d) \quad (4)$$

where a is the relative surface fraction of the aligned (nematic) domains. It can further be shown that decrease of the order with the box size d is an exponential function⁶⁵

$$S_{2D}^{\text{align}}(d) = b + (1 - b)e^{-d/2\Lambda} \quad (5)$$

where the offset b accounts the finite size of the image and Λ is the characteristic length of the nematic order decay, in analogy to the relationship between the persistence length and tangent correlations along the contour. The component $S_{2D}^{\text{rand}}(d)$ can be evaluated from simulated images of randomly oriented fibrils that are generated using all relevant parameters from the original AFM images.

Applying this analysis to the tracked contours of β -lactoglobulin amyloid fibrils at the air–water interface after different adsorption times (Figure 8D–G), it is possible to quantify isotropic–nematic transition in a reliable and rigorous manner.⁶⁵

From 2D Projection Traces to 3D Structural Parameters. Much of the discussion above refers to objects with a 2D conformation because either adsorbed on a substrate (AFM, TEM) or with a pseudo-2D conformation when confined within a thin layer of suspension (Cryo-TEM), for which a 2D polymer physics analysis provides an accurate description. It is however interesting to extract also information for objects, which span a truly 3D conformation, but for which only their projection traces on a plane are accessible. This is the case, for example for fibrils visualized by confocal microscopy or fluorescence-based stochastic optical reconstruction microscopy (STORM). In order to access the persistence length of fibrillar structures distributed randomly in 3D, one should carefully consider how the statistical properties of WLC projections change from a 3D space into a 2D plane. When a chain is projected onto the xy plane from the z direction, the mean-square of the projected end-to-end distance modifies into $\langle R^2 \rangle_{\text{proj}} = (2/3)\langle R^2 \rangle_{3D}$.⁹ The observable projected internal contour length l_{proj} also differs from the actual one, l_{3D} . To understand how, let l_{3D} be decomposed in N segments of identical length ΔS so that $l_{3D} = N\Delta S$; similarly, $l_{\text{proj}} = \sum_i^N \Delta S_{\text{proj}}^i$ where ΔS_{proj}^i are the individual projections of the various ΔS^i on the projection plane. Then we can rewrite $l_{\text{proj}} = N[(1/N)\sum_i^N \Delta S_{\text{proj}}^i]$, or $l_{\text{proj}} = N\langle \Delta S_{\text{proj}} \rangle$, where $\langle \Delta S_{\text{proj}} \rangle$ is

nothing else than the average projection of ΔS^i on the projection plane. By averaging over all possible directions of segments ΔS placed at an angle θ from z , we derive $\langle \Delta S_{\text{proj}} \rangle = (2/\pi) \int_0^{\pi/2} \Delta S \sin \theta \, d\theta = (2/\pi)\Delta S$, and then $l_{\text{proj}} = N\langle \Delta S_{\text{proj}} \rangle = N(2/\pi)\Delta S = (2/\pi)l_{3D}$. By applying the WLC model in 3D expressing $\langle R^2 \rangle_{3D}$ vs l_{3D} for the actual fibril, $\langle R^2 \rangle_{3D} = 2\lambda[l_{3D} - \lambda(1 - e^{-l_{3D}/\lambda})]$, but using only the observable quantities $\langle R^2 \rangle_{\text{proj}}$ and l_{proj} , one gets $\langle R^2 \rangle_{\text{proj}} = (2\lambda/3)[\pi l_{\text{proj}} - 2\lambda(1 - e^{-\pi l_{\text{proj}}/2\lambda})]$, where λ is the real persistence length of WLC in 3D. This illustrates well that in the case of randomly oriented WLC fibrils in 3D their structural parameters, such as the persistence length, are easily accessible through a statistical analysis of their projections traces.

Beside the persistence length, other important structural parameters can be extracted from the projection traces. For example, because fibrils are rather rigid, $\langle R^2 \rangle_{3D} \sim l_{3D}^\nu$, with the scaling exponent ν typically ranging between the very flexible random walk value, $\nu = 1/2$, and the fully rigid chain value, $\nu = 1$, the fractal exponent $m = 1/\nu$ ranges between 1 and 2 and is thus preserved upon projection on a surface of Euclidean dimension two,⁸⁹ implying that the scaling and fractal exponents for the real fibril and its projection are identical and thus accessible via the observable $\langle R^2 \rangle_{\text{proj}} \sim l_{\text{proj}}^\nu$.

OUTLOOK AND CONCLUSIONS

Given the paramount significance that elongated, semiflexible colloidal objects and biological macromolecules embody in medicine, biology, materials science, and nanotechnology, there is an increasing urgent need to develop semiautomated procedures for the processing of a large amount of microscopy images and generating statistical significant output on the structural features and molecular conformations of these filamentous objects. The newly developed FiberApp open source code is deeply rooted in statistical polymer physics and is meant to bridge the gap currently existing between the high-resolution imaging process of fiber-like, filamentous, and macromolecular objects and their structural analysis at the single molecule level. We anticipate that this code will serve a vast community of scientists working on very diverse disciplines and may pave the way to new approaches to the study of biological and synthetic polymers such as microtubules, actin, amyloid fibrils, collagen, nanocellulose, silk fibroin, and the many types of supramolecular fibers found in current applied and fundamental sciences.

ASSOCIATED CONTENT

Supporting Information

Technical details of processing algorithms as well as specifications of A* pathfinding and active contour models based tracking procedures; FiberApp software. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Notes

The authors declare no competing financial interest.

Biographies



Ivan Usov was born in 1988 in Cheboksary, Russia. In 2011, under the guidance of Prof. A. Muzafarov, he received his Master's degree with Honors on specialty Physics of Condensed Matter from Moscow State University, Russia. He then joined the group of Prof. R. Mezzenga at ETH Zurich to pursue a PhD degree on statistical analysis of fibrous objects and filaments. His research interests include quantitative analysis, chirality, and mechanical properties of natural fibrillar nanostructures.



Raffaele Mezzenga received his PhD in the field of polymer physics from EPFL Lausanne. He was a postdoc at University of California, Santa Barbara, in 2001–2002. He then moved to Nestlé Research Center to work on the self-assembly of surfactants, natural amphiphiles, and liquid crystals. In 2005 he was hired as Associate Professor in the Physics Department of the University of Fribourg, and he then joined ETH Zurich on 2009 as Full Professor. His research focuses on the fundamental understanding of self-assembly processes in polymers, liquid crystals, food, and biological colloidal systems. Prof. Mezzenga has coauthored more than 200 publications and is recipient of several international awards, among which are the 2004 Swiss Science National Foundation Professorship, the 2011 Young Scientist Research Award (AOCS), the 2011 John Dillon Medal (APS), and the 2013 Biomacromolecules/Macromolecules Young Investigator Award (ACS).

ACKNOWLEDGMENTS

Authors acknowledge support from the Swiss National Science Foundation (SNF 2-77002-11) and are grateful to Larissa Schefer and Sophia Jordens for inspiring discussions and valuable assistance in software testing. The critical reading of the manuscript by Kathleen Smith and Ilya Savchenko is kindly acknowledged.

REFERENCES

- (1) Meijering, E. *Cytometry, Part A* **2010**, *77*, 693–704.
- (2) Fan, J.; Zhou, X.; Dy, J. G.; Zhang, Y.; Wong, S. T. C. *Neuroinformatics* **2009**, *7*, 113–130.
- (3) Yu, W.; Lee, H. K.; Hariharan, S.; Bu, W.; Ahmed, S. *Cytometry, Part A* **2009**, *75*, 289–297.
- (4) Leandro, J. J.; Cesar, R. M., Jr.; Costa Lda, F. J. *Neurosci. Methods* **2009**, *177*, 497–509.
- (5) Machens, C. K. *Science* **2012**, *338*, 1156–1157.
- (6) Fujiwara, I.; Vavylonis, D.; Pollard, T. D. *Proc. Natl. Acad. Sci. U. S. A.* **2007**, *104*, 8827–8832.
- (7) Kuhn, J. R.; Pollard, T. D. *Biophys. J.* **2005**, *88*, 1387–1402.
- (8) Fujiwara, I.; Takahashi, S.; Tadakuma, H.; Funatsu, T.; Ishiwata, S. *Nat. Cell Biol.* **2002**, *4*, 666–673.
- (9) Rivetti, C.; Guthold, M.; Bustamante, C. *J. Mol. Biol.* **1996**, *264*, 919–932.
- (10) Valle, F.; Favre, M.; De Los Rios, P.; Rosa, A.; Dietler, G. *Phys. Rev. Lett.* **2005**, *95*, 158105.
- (11) Faas, F. G. A.; Rieger, B.; van Vliet, L. J.; Cherny, D. I. *Biophys. J.* **2009**, *97*, 1148–1157.
- (12) Witz, G.; Rechendorff, K.; Adamcik, J.; Dietler, G. *Phys. Rev. Lett.* **2008**, *101*, 148103.
- (13) Schleeper, M.; VandenAkker, C. C.; Deckert-Gaudig, T.; Deckert, V.; Velikov, K. P.; Koenderink, G.; Bonn, M. *Polymer* **2013**, *54*, 2473–2488.
- (14) Knowles, T. P.; Fitzpatrick, A. W.; Meehan, S.; Mott, H. R.; Vendruscolo, M.; Dobson, C. M.; Welland, M. E. *Science* **2007**, *318*, 1900–1903.
- (15) Volpatti, L. R.; Knowles, T. P. J. *J. Polym. Sci., Part B: Polym. Phys.* **2014**, *52*, 281–292.
- (16) Adamcik, J.; Lara, C.; Usov, I.; Jeong, J. S.; Ruggeri, F. S.; Dietler, G.; Lashuel, H. A.; Hamley, I. W.; Mezzenga, R. *Nanoscale* **2012**, *4*, 4426–4429.
- (17) Cherny, I.; Gazit, E. *Angew. Chem., Int. Ed.* **2008**, *47*, 4062–4069.
- (18) Knowles, T. P. J.; Buehler, M. J. *Nat. Nanotechnol.* **2011**, *6*, 469–479.
- (19) Dobson, C. M. *Nature* **2003**, *426*, 884–890.
- (20) Selkoe, D. J. *Nature* **2003**, *426*, 900–904.
- (21) Wang, J. C.; Turner, M. S.; Agarwal, G.; Kwong, S.; Josephs, R.; Ferrone, F. A.; Briehl, R. W. *J. Mol. Biol.* **2002**, *315*, 601–612.
- (22) Yoon, G.; Kab Kim, Y.; Eom, K.; Na, S. *Appl. Phys. Lett.* **2013**, *102*, 011914.
- (23) Rust, M. J.; Bates, M.; Zhuang, X. *Nat. Methods* **2006**, *3*, 793–795.
- (24) Albertazzi, L.; van der Zwaag, D.; Leenders, C. M. A.; Fitzner, R.; van der Hofstad, R. W.; Meijer, E. W. *Science* **2014**, *344*, 491–495.
- (25) Smith, M. B.; Li, H.; Shen, T.; Huang, X.; Yusuf, E.; Vavylonis, D. *Cytoskeleton* **2010**, *67*, 693–705.
- (26) Vasilkoski, Z.; Stepanyants, A. *J. Neurosci. Methods* **2009**, *178*, 197–204.
- (27) Hall, D. *Anal. Biochem.* **2012**, *421*, 262–277.
- (28) Haralick, R. M.; Watson, L. T.; Laffey, T. J. *Int. J. Rob. Res.* **1983**, *2*, 50–72.
- (29) Kass, M.; Witkin, A.; Terzopoulos, D. *Int. J. Comput. Vis.* **1988**, *1*, 321–331.
- (30) Steger, C. *IEEE Trans. Pattern Anal. Mach. Intell.* **1998**, *20*, 113–125.
- (31) Jacob, M.; Unser, M. *Pattern Anal. Mach. Intell.* **2004**, *26*, 1007–1019.
- (32) Sun, C.; Vallotton, P. J. *Microsc.* **2009**, *234*, 147–157.
- (33) Frangi, A. F.; Niessen, W. J.; Hoogeveen, R. M.; van Walsum, T.; Viergever, M. A. *IEEE Trans. Med. Imaging* **1999**, *18*, 946–956.
- (34) Lorigo, L. M.; Faugeras, O. D.; Grimson, W. E.; Keriven, R.; Kikinis, R.; Nabavi, A.; Westin, C. F. *Med. Image Anal.* **2001**, *5*, 195–206.
- (35) Sato, Y.; Nakajima, S.; Shiraga, N.; Atsumi, H.; Yoshida, S.; Koller, T.; Gerig, G.; Kikinis, R. *Med. Image Anal.* **1998**, *2*, 143–168.

- (36) Yuan, X.; Trachtenberg, J. T.; Potter, S. M.; Roysam, B. *Neuroinformatics* **2009**, *7*, 213–232.
- (37) Broser, P. J.; Erdogan, S.; Grinevich, V.; Osten, P.; Sakmann, B.; Wallace, D. J. *J. Neurosci. Methods* **2008**, *169*, 43–54.
- (38) Narro, M. L.; Yang, F.; Kraft, R.; Wenk, C.; Efrat, A.; Restifo, L. *Brain Res.* **2007**, *1138*, 57–75.
- (39) Pool, M.; Thiemann, J.; Bar-Or, A.; Fournier, A. E. *J. Neurosci. Methods* **2008**, *168*, 134–139.
- (40) Janoos, F.; Mosaliganti, K.; Xu, X.; Machiraju, R.; Huang, K.; Wong, S. T. C. *Med. Image Anal.* **2009**, *13*, 167–179.
- (41) Hart, P.; Nilsson, N.; Raphael, B. *IEEE Trans. Syst. Sci. Cybern.* **1968**, *4*, 100–107.
- (42) Longair, M. H.; Baker, D. A.; Armstrong, J. D. *Bioinformatics* **2011**, *27*, 2453–2454.
- (43) Meijering, E.; Jacob, M.; Sarria, J.-C. F.; Steiner, P.; Hirling, H.; Unser, M. *Cytometry, Part A* **2004**, *58*, 167–176.
- (44) Schmitt, S.; Evers, J. F.; Duch, C.; Scholz, M.; Obermayer, K. *Neuroimage* **2004**, *23*, 1283–1298.
- (45) Rodriguez, A.; Ehlenberger, D. B.; Hof, P. R.; Wearne, S. L. *J. Neurosci. Methods* **2009**, *184*, 169–175.
- (46) Nisslert, R.; Kvarnström, M.; Lorén, N.; Nydén, M.; Rudemo, M. *J. Microsc.* **2007**, *225*, 10–21.
- (47) Altendorf, H.; Decencièrre, E.; Jeulin, D.; De sa Peixoto, P.; Deniset-Besseau, A.; Angelini, E.; Mosser, G.; Schanne-Klein, M.-C. *J. Microsc.* **2012**, *247*, 161–175.
- (48) Wu, J.; Rajwa, B.; Filmer, D. L.; Hoffmann, C. M.; Yuan, B.; Chiang, C.; Sturgis, J.; Robinson, J. P. *J. Microsc.* **2003**, *210*, 158–165.
- (49) Stein, A. M.; Vader, D. A.; Jawerth, L. M.; Weitz, D. A.; Sander, L. M. *J. Microsc.* **2008**, *232*, 463–475.
- (50) Brangwynne, C. P.; Koenderink, G. H.; Barry, E.; Dogic, Z.; MacKintosh, F. C.; Weitz, D. A. *Biophys. J.* **2007**, *93*, 346–359.
- (51) Hadjidelmetriou, S.; Toomre, D.; Duncan, J. *Med. Image Anal.* **2008**, *12*, 689–702.
- (52) Wiggins, P. A.; van der Heijden, T.; Moreno-Herrero, F.; Spakowitz, A.; Phillips, R.; Widom, J.; Dekker, C.; Nelson, P. C. *Nat. Nanotechnol.* **2006**, *1*, 137–141.
- (53) Moukhtar, J.; Faivre-Moskalenko, C.; Milani, P.; Audit, B.; Vaillant, C.; Fontaine, E.; Mongelard, F.; Lavorel, G.; St-Jean, P.; Bouvet, P.; Argoul, F.; Arneodo, A. *J. Phys. Chem. B* **2010**, *114*, 5125–5143.
- (54) Adamcik, J.; Jung, J.-M.; Flakowski, J.; De Los Rios, P.; Dietler, G.; Mezzenga, R. *Nat. Nanotechnol.* **2010**, *5*, 423–428.
- (55) Smith, J. F.; Knowles, T. P. J.; Dobson, C. M.; Macphee, C. E.; Welland, M. E. *Proc. Natl. Acad. Sci. U. S. A.* **2006**, *103*, 15806–15811.
- (56) Witz, G.; Rechendorff, K.; Adamcik, J.; Dietler, G. *Phys. Rev. Lett.* **2011**, *106*, 248301.
- (57) Schefer, L.; Adamcik, J.; Mezzenga, R. *Angew. Chem., Int. Ed.* **2014**, *53*, 5376–5379.
- (58) Drube, F.; Alim, K.; Witz, G.; Dietler, G.; Frey, E. *Nano Lett.* **2010**, *10*, 1445–1449.
- (59) Sakaue, T.; Witz, G.; Dietler, G.; Wada, H. *EPL* **2010**, *91*, 68002.
- (60) Timoshenko, E. G.; Kuznetsov, Y. A.; Connolly, R. *J. Chem. Phys.* **2002**, *116*, 3905.
- (61) Peng, H.; Ruan, Z.; Long, F.; Simpson, J. H.; Myers, E. W. *Nat. Biotechnol.* **2010**, *28*, 348–353.
- (62) Mikhaylov, A.; Sekatskii, S. K.; Dietler, G. *J. Adv. Microsc. Res.* **2013**, *8*, 241–245.
- (63) Usov, I.; Adamcik, J.; Mezzenga, R. *ACS Nano* **2013**, *7*, 10465–10474.
- (64) Usov, I.; Adamcik, J.; Mezzenga, R. *Faraday Discuss.* **2013**, *166*, 151–162.
- (65) Jordens, S.; Isa, L.; Usov, I.; Mezzenga, R. *Nat. Commun.* **2013**, *4*, 1917.
- (66) Jordens, S.; Riley, E. E.; Usov, I.; Isa, L.; Olmsted, P. D.; Mezzenga, R. *ACS Nano* **2014**, *8*, 11071–11079.
- (67) Li, C.; Mezzenga, R. *Langmuir* **2012**, *28*, 10142–10146.
- (68) Rivetti, C.; Walker, C.; Bustamante, C. *J. Mol. Biol.* **1998**, *280*, 41–59.
- (69) Lara, C.; Usov, I.; Adamcik, J.; Mezzenga, R. *Phys. Rev. Lett.* **2011**, *107*, 238101.
- (70) Bolisetty, S.; Vallooran, J. J.; Adamcik, J.; Mezzenga, R. *ACS Nano* **2013**, *7*, 6146–6155.
- (71) Usov, I.; Nyström, G.; Adamcik, J.; Handschin, S.; Bergström, L.; Mezzenga, R., submitted.
- (72) Lu, Y.; Weers, B. D.; Stellwagen, N. C. *Biophys. J.* **2003**, *85*, 409–415.
- (73) Dame, R. T.; van Mameren, J.; Luijsterburg, M. S.; Mysiak, M. E.; Janičević, A.; Pazdzior, G.; van der Vliet, P. C.; Wyman, C.; Wuite, G. J. L. *Nucleic Acids Res.* **2005**, *33*, e68.
- (74) Saha, A.; Adamcik, J.; Bolisetty, S.; Handschin, S.; Mezzenga, R., submitted.
- (75) Lara, C.; Reynolds, N. P.; Berryman, J. T.; Xu, A.; Zhang, A.; Mezzenga, R. *J. Am. Chem. Soc.* **2014**, *136*, 4732–4739.
- (76) McGeoch, C. C. *A Guide to Experimental Algorithmics*; Cambridge University Press: Cambridge, UK, 2012; p 171.
- (77) Usov, I.; Mezzenga, R. *ACS Nano* **2014**, *8*, 11035–11041.
- (78) Rubinstein, M.; Colby, R. H. *Polymer Physics*; Oxford University Press: New York, 2003; pp 54–60.
- (79) Wiggins, P. A.; Nelson, P. C. *Phys. Rev. E: Stat., Nonlinear, Soft Matter Phys.* **2006**, *73*, 031906.
- (80) Hsu, H.-P.; Paul, W.; Binder, K. *Macromolecules* **2010**, *43*, 3094–3102.
- (81) Doi, M.; Edwards, S. F. *The Theory of Polymer Dynamics*; Oxford University Press: New York, 1986; p 317.
- (82) Manning, G. S. *Phys. Rev. A* **1986**, *34*, 668–670.
- (83) De Gennes, P. G. *Scaling Concepts in Polymer Physics*; Cornell University Press: Ithaca, NY, 1979.
- (84) Mücke, N.; Klenin, K.; Kirmse, R.; Bussiek, M.; Herrmann, H.; Hafner, M.; Langowski, J. *PLoS One* **2009**, *4*, e7756.
- (85) Rappaport, S. M.; Medalion, S.; Rabin, Y. arXiv:0801.3183, 2008.
- (86) Villarrubia, J. S. *Surf. Sci.* **1994**, *321*, 287–300.
- (87) Villarrubia, J. S. *J. Res. Natl. Inst. Stand. Technol.* **1997**, *102*, 425–454.
- (88) Lamour, G.; Yip, C. K.; Li, H.; Gsponer, J. *ACS Nano* **2014**, *8*, 3851–3861.
- (89) Mandelbrot, B. B. *The Fractal Geometry of Nature*; W.H. Freeman and Company: San Francisco, 1982.